## ФИЗИКА

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## TEMPERATURE DEPENDENCE OF THE FREQUENCY OF TWO-ELECTRON EXCHANGE BETWEEN IMPURITY NEGATIVE-U TIN-CENTERS IN LEAD SULFIDE

A Fast two\_electron exchange between neutral  $Sn^{2+}$  and doubly ionized  $Sn^{4+}$  impurity negative\_U tin centers in partially compensated  $Pb_{0.98}Sn_{0.02}Na_{0.01}Tl_{0.01}S$  solid solutions has been found by emission Mössbauer spectroscopy on  $^{119mm}Sn(^{119m}Sn)$  isotope; the lifetime of the  $Sn^{2+}$  and  $Sn^{4+}$  states changes from ~  $6 \times 10^{-4}$  to ~  $8 \times 10^{-9}$  s with a change in temperature from 295 to 900 K.

Keywords: Mössbauer spectroscopy, two electron exchange, U-centers.

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## ТЕМПЕРАТУРНАЯ ЗАВИСИМОСТЬ ЧАСТОТЫ ДВУХЭЛЕКТРОННОГО ОБМЕНА МЕЖДУ ПРИМЕСНЫМИ U-МИНУС-ЦЕНТРАМИ ОЛОВА В СУЛЬФИДЕ СВИНЦА

Методом эмиссионной мессбауэровской спектроскопии на изотопе  $^{119mm}$ Sn( $^{119m}$ Sn) обнаружен процесс быстрого двухэлектронного обмена между нейтральными Sn<sup>2+</sup> и двукратно ионизованными Sn<sup>4+</sup> примесными U-минус-центрами олова в частично компенсированных твердых растворах  $Pb_{0.98}$ Sn<sub>0.02</sub>Na<sub>0.01</sub>Tl<sub>0.01</sub>S, причем время жизни состояний Sn<sup>2+</sup> и Sn<sup>4+</sup> меняется от ~ 6 · 10<sup>-4</sup> с до ~ 8 · 10<sup>-9</sup> с при изменении температуры от 295 до 900 K.

**Ключевые слова:** мессбауэровская спектроскопия, двухэлектронный обмен, U-минусцентры.

The notion of negative-U centers was used for the first time by Anderson [1] to explain the electric and magnetic properties of glassy semiconductors. Mössbauer spectroscopy on <sup>119</sup>Sn isotope was found to be the most efficient method for identifying negative-U centers in glassy and crystalline semiconductors, as well as superconductors [2, 3]. In particular, it was shown that impurity tin atoms in lead chalcogenides (PbS and PbSe) are located in regular sites of cationic sublattice and form two-electron donor centers with negative correlation energy; the parameters of the microscopic model of negative-U centers were also determined. However, the question of the existence of two-electron exchange between neutral and ionized negative-U centers in semiconductors remains open. In this article, we report on the uncovering of this process and determining the temperature dependence of the frequency of two-electron exchange between neutral and ionized tin centers in partially compensated solid solutions based on lead sulfide by emission Mössbauer spectroscopy on <sup>119mm</sup>Sn(<sup>119m</sup>Sn) isotope in the temperature range of 80–900 K.

 $Pb_{0.98}Sn_{0.02}Na_{0.01}Tl_{0.01}S$  solid solutions were prepared by alloying the initial components of semiconductor purity grade in evacuated quartz ampoules with subsequent annealing ingots and then pressed powders at 650°C for 120 h. All samples were single-phase and had a NaCl-type structure. The samples had p-type conductivity and were nondegenerate (the hole concentration varied from ~ 5 × 10<sup>13</sup> at 80 K to ~ 10<sup>17</sup> cm<sup>-3</sup> at 295 K).  $Pb_{0.98}^{119mm}Sn_{0.02}Na_{0.01}Tl_{0.01}S$  Mössbauer sources were prepared using a <sup>118</sup>Sn preparation irradiated by a neutron flux of ~ 10<sup>15</sup> cm<sup>-2</sup> s<sup>-1</sup> for 6 months.

In figure 1 Emission Mössbauer spectra of  ${}^{119\text{mm}}\text{Sn}({}^{119\text{m}}\text{Sn})$  isotope in Pb<sub>0.98</sub>Sn<sub>0.02</sub>Na<sub>0.01</sub>Tl<sub>0.01</sub>S solid solutions recorded at 80 and 295 K are shown. They are superpositions of two lines of instrumental width ( $G_{sn^{2+}} = G_{sn^{4+}} = 0.80(1)$  mm/sm), isomeric shifts  $\delta$  of which correspond to the Sn<sup>2+</sup> ( $\delta_{sn^{2+}} = 3.69(1)$  mm/s) and Sn<sup>4+</sup> ( $\delta_{sn^{4+}} = 1.26(1)$  mm/s) centers.

Fig. 1. Emission Mössbauer spectra of  $^{119mm}$ Sn( $^{119m}$ Sn) isotope in Pb<sub>0.98</sub>Sn<sub>0.02</sub>Na<sub>0.01</sub>Tl<sub>0.01</sub>S solid solutions recorded by a resonant detector at source temperatures of 80, 295, 400, 500, and 900 K. The lines corresponding to the Sn<sup>2+</sup> and Sn<sup>4+</sup> centers are shown. The spectrum recorded at 900 K corresponds to the averaged state of tin centers



An increase in temperature from 80 to 295 K is accompanied by a decrease in relative intensity R of the  $\text{Sn}^{2+}$  line (R = 0.49(1) at 80 K and R = 0.41(2) at 295 K), which can be explained by a sharper temperature dependence of the fraction of recoilless processes for divalent tin compounds, as compared with four-valence tin compounds. The  $\text{Sn}^{2+}$  and  $\text{Sn}^{4+}$  lines are somewhat broadened ( $G_{sn^{2+}} = G_{sn^{4+}} = 0.90(1)$  mm/s) and located closer to each other ( $\delta_{sn^{2+}} = 3.62(1)$  mm/s) and Sn<sup>4+</sup> ( $\delta_{sn^{4+}} = 1.36(1)$  mm/s).

At 500 K (figure 1), the relative intensity of the Sn<sup>2+</sup> line further decreases (R = 0.31(3)) and the Sn<sup>2+</sup> and Sn<sup>4+</sup> lines are significantly broadened ( $G_{Sn^{2+}} = G_{Sn^{4+}} = 1.6(1)$  mm/s) and even closer to each other ( $\delta_{Sn^{2+}} = 1.5(1)$  mm/s and  $\delta_{Sn^{4+}} = 3.2(1)$  mm/s). Finally, at 900 K (figure 1), the lines are merged into one broadened line (G = 2.4(1) mm/s), the isomeric shift of which ( $\delta = 2.0(2)$  mm/s) is intermediate between those for the Sn<sup>2+</sup> and Sn<sup>4+</sup> centers.

The spectra in figure 1 illustrate a typical process of electron exchange between two  $\text{Sn}^{2+}$  and  $\text{Sn}^{4+}$  states; the frequency of electron exchange increases with an increase in temperature. The spectrum recorded at 900 K corresponds to the averaged state of tin centers, which arises due to the fast electron exchange between  $\text{Sn}^{2+}$  and  $\text{Sn}^{4+}$ . The absence of the intermediate charge state of tin centers ( $\text{Sn}^{3+}$ ) in the spectra indicates that the exchange occurs via simultaneous transfer of two electrons.

To determine the frequency of electron exchange, the experimental spectra were processed by the least-squares method on the assumption that the shape of the spectral line is determined by the relation

$$W(V) = -\frac{AC + BD}{C^2 + D^2},$$

where:

$$\begin{split} A &= J_{Sn^{2+}} \left( G_{Sn^{4+}} + \tau_{Sn^{4+}}^{-1} + \tau_{Sn^{2+}}^{-1} \right), \\ B &= J_{Sn^{2+}} \left( \delta_{Sn^{4+}} - V \right) + J_{Sn^{4+}} \left( \delta_{Sn^{2+}} - V \right), \\ C &= \left( G_{Sn^{2+}} + \tau_{Sn^{2+}}^{-1} \right) \left( G_{Sn^{4+}} + \tau_{Sn^{4+}}^{-1} \right) - \left( \delta_{Sn^{2+}} - V \right) \left( \delta_{Sn^{4+}} - V \right) - \tau_{Sn^{2+}}^{-1} \tau_{Sn^{4+}}^{-1}, \\ D &= \left( \delta_{Sn^{2+}} - V \right) \left( G_{Sn^{4+}} + \tau_{Sn^{4+}}^{-1} \right) + \left( \delta_{Sn^{4+}} - V \right) \left( G_{Sn^{2+}} + \tau_{Sn^{2+}}^{-1} \right), \end{split}$$

 $J_{Sn^{2+}}$  and  $J_{Sn^{4+}}$  — are the amplitudes of the Sn<sup>2+</sup> and Sn<sup>4+</sup> lines, respectively; and  $\tau_{Sn^{2+}}$  and  $\tau_{Sn^{2+}}$  are the lifetimes of the Sn<sup>2+</sup> and Sn<sup>4+</sup> centers, respectively (we assume that  $\tau_{Sn^{2+}} = \tau_{Sn^{4+}} = \tau$  because of the lack of experimental parameters).

The temperature dependence of the frequency of electron exchange  $v = \tau^{-1}$  between the centers of the Sn<sup>2+</sup> and Sn<sup>4+</sup> is shown in figure 2.



The activation energy of exchange is 0.11(2) eV. This corresponds to the distance of the Fermi level from the valence band to the hole partially-compensated solid solutions  $Pb_{1-x-1}$ 

 $_{y}$ Sn<sub>x</sub>(Na,Tl)<sub>y</sub>S [2], and indicates that the electron exchange between the centers of Sn<sup>2+</sup> and Sn<sup>4+</sup> is implemented using state of the valence band. In favor of such a mechanism is evidenced by the fact that the exchange is observed at low concentrations of tin, when the process can not be a direct exchange of electrons between the centers of the tin.

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# CORRELATIONS OF THE <sup>63</sup>Cu NMR DATA WITH THE <sup>67</sup>Cu (<sup>67</sup>Zn) AND THE <sup>61</sup>Cu (<sup>61</sup>Ni) EMISSION MÖSSBAUER DATA FOR CERAMIC SUPERCONDUCTORS

A linear correlation between the quadrupole coupling constant  $C_{Cu}$  measured by the <sup>63</sup>Cu NMR technique on the one hand and the quadrupole coupling constants  $C_{Zn}$  and  $C_{Ni}$  measured by the <sup>67</sup>Cu(<sup>67</sup>Zn) and <sup>61</sup>Cu(<sup>61</sup>Ni) emission Mössbauer spectroscopy on the other hand has been found for ceramic superconductors.

Keywords: Mössbauer spectroscopy, NMR, electric field gradient.

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### КОРРЕЛЯЦИОННЫЕ СООТНОШЕНИЯ МЕЖДУ ДАННЫМИ ЯМР <sup>63</sup>Си И ЭМИССИОННОЙ МЕССБАУЭРОВСКОЙ СПЕКТРОСКОПИИ <sup>67</sup>Си(<sup>67</sup>Zn) И <sup>61</sup>Сu(<sup>61</sup>Ni) ДЛЯ КЕРАМИЧЕСКИХ СВЕРХПРОВОДНИКОВ

Установлена линейная корреляция между постоянной квадрупольного расщепления  $C_{Cu}$ , измеренной методом ЯМР<sup>63</sup>Си и постоянными квадрупольного расщепления  $C_{Zn}$  и  $C_{Ni}$ , измеренные методом эмиссионной мессбауэровской спектроскопии<sup>67</sup>Си(<sup>67</sup>Zn) и <sup>61</sup>Си(<sup>61</sup>Ni) для керамических сверхпроводников.

Ключевые слова: мессбауэровская спектроскопия, ЯКР, градиент электрического поля.

#### 1. Introduction

One of the main problems in physics of high-temperature (high- $T_c$ ) superconductors is the determination of the spatial distribution of electronic defects in the lattices of copper metal oxides. A potentially effective method to solve this problem is to compare the experimentally determined and calculated parameters of the electric field gradient (EFG) tensor for specific lattice sites [2]. Copper sites are of utmost interest in such work because these atoms are found in the